

Band gap engineering of anatase TiO₂ by ambipolar doping: A first principles study

S. Muthukrishnan^{a,b}, R. Vidya^{a,b,*}, Anja Olafsen Sjøstad^c

^a Computational Laboratory for Multifunctional Materials (CoLaMM), Department of Medical Physics, Anna University, Sardar Patel Road, Guindy, Chennai, 600025, India

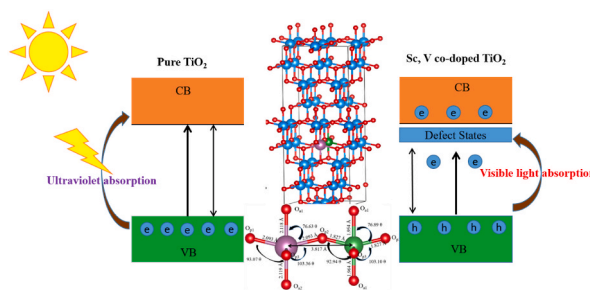
^b Centre for Materials Informatics (C-MaIn), Sir. C.V. Raman Science Block, Anna University, Sardar Patel Road, Guindy, Chennai, 600 025, India

^c Centre Materials Science and Nanotechnology, Department of Chemistry, University of Oslo, P.O. Box: 1033, Blindern, N-0315, Oslo, Norway

HIGHLIGHTS

- Engineering the band gap and by co-doping with Sc and V impurities in TiO₂.
- Calculation of defect energetics and bonding characteristics by the impurities on anatase TiO₂.
- Exploring the effect of Jahn-Teller distortion by mono-doping and co-doping in Anatase TiO₂.
- Enhanced optical absorption in the visible region of the solar spectrum.

GRAPHICAL ABSTRACT



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ABSTRACT

Pure Anatase TiO₂ is limited to photocatalytic activity only in the UV region of solar energy, (i.e. $E_g = 3.2$ eV and $\lambda < 388$ nm), which utilizes only 5% of the solar spectrum. Adding external impurities to Anatase TiO₂ facilitates absorption of the visible spectrum by reducing the band gap of TiO₂. Hence, we investigated the effect of Sc and V mono-doping and co-doping on the pure anatase TiO₂ by Density Functional Theory (DFT). To elucidate the effect of doping on the electronic structure, the accurate band structure and density of states are derived using the hybrid functional methodology (HSE06). The defect formation energy of the dopants in their different charge states is calculated using the gradient-corrected density functionals using the PBE method. The bonding characteristics were analyzed using the charge density and electron-localization function (ELF) plots, which indicate that the impurities cause a local lattice distortion, which further influences the band features of TiO₂. It is observed that the formation of additional energy states by the external impurities below the conduction band minimum, reduces the band gap of the material significantly. The optical absorption analysis shows that Sc and V co-doping not only improves the absorption in the UV region, but also reaches redshift absorption. Hence the ambipolar co-doping by Sc and V to the pure anatase TiO₂ makes it a better for visible light active photocatalyst.

* Corresponding author. Centre for Materials Informatics (C-MaIn), Sir. C.V. Raman Science Block, Anna University, Sardar Patel Road, Guindy, Chennai, 600 025, India.

E-mail addresses: vidyar@annauniv.edu, vidyapl.ravi@gmail.com (R. Vidya).

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1. Introduction

TiO₂ is considered as the most promising semiconductor material, with enormous applications in the field of photo catalysts for pollution reduction, sensors, solar cell, water splitting, hydrogen production, and adsorption of the harmful organic molecules [1]. Naturally occurring TiO₂ exists in three polymorphs known as Anatase, Rutile and Brookite [2], with a bandgap (E_g) value of 3.2, 3.02, and 3.3 eV, respectively [3]. In 1972, Fujishima and Honda found that photo electrochemical splitting of water (H₂O into H₂ and O₂) can be achieved by irradiating the TiO₂ surface with UV light [4]. Due to the large bandgap of TiO₂, it can be operated only in the UV region as a photo catalyst. However, to carry out the redox reaction, charge carriers must be generated in the visible region of the solar spectrum [5]. Therefore, TiO₂ is activated at a wavelength of 388 nm (corresponding to the bandgap <3.2 eV) for redox reaction, which is only a small fraction of the entire solar spectrum [6]. TiO₂ can be made more functional in absorbing visible light by doping other elements to them [7]. In order to achieve the visible light-driven photo catalysts, the following conditions are necessary, (i) reducing the bandgap energy less than 3 eV, (ii) making appropriate potentials for CB and VB edges, and (iii) increasing the mobility of electron and hole charge carriers [8]. Doping the TiO₂ modifies its bandgap value by increasing the charge carriers' separation and induces more light absorption in the visible region [9]. N-doped TiO₂ shows red-shifted absorption toward the visible light, because defect states are induced above the valence band maximum by the dopant atoms [10]. The reason behind the visible light absorption of TiO₂ may be due to the bandgap narrowing and reduction in the recombination rate of electron and hole charge carriers [11]. The key mechanism which enhances the photocatalytic activity is the formation of additional energy levels by the dopants inside the bandgap [12]. The formation of defect energy level not only enhances the absorption of photons in the UV region but also in the visible region [13].

The formation of additional energy levels in the forbidden gap may act as an electron trap, which eventually reduces the recombination of electron and hole pairs [14]. Reducing the charge carrier recombination rate is the vital step for increasing the photocatalytic activity [15]. Co-doping methods enhance the photocatalytic activity by overcoming the drawbacks of mono doping, where defect level may act as the trap sites for charge carriers. Xiong et al. studied the Sc and S co-doped anatase TiO₂, and found that the synergic effect induced by co-doping increases light absorption towards the visible region [16,17]. To maintain the charge neutrality, ambipolar doping is considered. Lin et al. studied anatase TiO₂ by first principles theory, and they found that co-doping with N and Si, decreases the bandgap to 3.01 eV. The optical absorption is also improved by widening the valence band to 5.3 eV because of the Si 3p orbitals [18]. Yanming et al. studied the optical properties of C and B co-doped anatase and rutile TiO₂, by density functional theory. The generation of additional energy states reduces its bandgap by 0.8 eV, which leads to red shift optical absorption on TiO₂. Co-doping generates the hybridized energy states at the bandgap by the mixing of 2p orbitals of C and B atoms. They also found that the visible light absorption is improved by increasing the C and B concentration in TiO₂ [19]. Khan et al. studied the W and Ag co-doped anatase TiO₂ through first principle study. They found that W and Ag co-doping improves the separation of photo generated charge carriers, and reduces its bandgap. Consequently the photocatalytic visible light absorbance was improved compared to the pure TiO₂ [20]. Xuechao et al. studied the electronic and optical properties of Cr and C co-doped anatase TiO₂. The lattice distortion caused by the co-doping makes a local internal electric field, thereby enhancing the separation of the photo excited electron and hole charge carriers. It is interesting to note that, Cr and C co-doping not only reduces the bandgap, but also inhibits the recombination of electron and hole charge carriers [21]. Cristiana et al. studied the Cr and Sb co-doped TiO₂. In addition to Cr doping, an oxygen vacancy is required to induce Cr³⁺ state for electron transfer. By making Cr and Sb

co-doping, Sb donates one electron to the Cr impurity atoms without the need Oxygen vacancies [22]. The excess charges generated by the dopants can be compensated by co-doping. Kulbir et al. studied the charge compensated Nb and Rh co-doped rutile TiO₂, where the creation of oxygen vacancy is not necessary for charge compensation [23].

One of the aims of the current study is to enhance the visible-light absorption of anatase TiO₂ by reducing its band-gap by doping. Even though the band-gap is reduced by doping with different impurity elements, the defect levels introduced by them sometimes act as recombination centres. In order to ensure charge neutrality, co-doping with simultaneous electron-donor and electron-acceptor impurities was also attempted. However, the variation in the atomic size of the co-doped elements leads to lattice distortion which may influence the behaviour of the host system under irradiation. Such lattice distortion may be minimized by introducing co-dopants of similar ionic radii. Therefore, we have taken the neighbors of Ti, viz Sc and V that have somewhat similar ionic radii like Ti for modelling the co-doping effect. By adding Sc³⁺ and V⁵⁺ to pure TiO₂, the excess charges of the material can be compensated. Hence the overall charges present in the structure remain neutral, consequently satisfying the valence demands. Murzin et al. successfully synthesised rutile TiO₂ with co-doping of Sc and V, and studied its structural and electrical properties [24].

2. Computational details

The simulation work based on Density Functional Theory was carried out by Plane-wave based pseudopotential method as implemented in Vienna Abinitio Simulation Package (VASP) for the self-consistent field total energy calculation [25,26]. The projected augmented wave method (PAW) exchange-correlation functional was followed. Perdew, Burke and Ernzerhof (PBE) Pseudopotentials were used to treat the core, and valence electron interaction [27]. The electron wave function was expanded by the plane wave basis set with a kinetic energy cut-off of 500 eV. A 2 × 2 × 2, and 4 × 4 × 4 k-point mesh were used for the Brillouin zone integration, for structure relaxation and electronic structure calculations respectively. The energy convergence was reached with setting the threshold value as 10⁻⁶ eV/atom, and the ions in the supercell were relaxed until forces on each atom less than 1 meV/Å by using conjugate gradient algorithm. A 2 × 2 × 2 supercell of 96 atoms of the TiO₂ anatase phase was considered for the GGA-PBE based calculations. For obtaining the band gap value and band dispersion accurately, Hybrid-DFT (HSE06 functional) was used [28]. A 2 × 2 × 1 supercell of 48 atoms of TiO₂ anatase phase was considered for the HSE06 based calculations. For the visualization of the crystal structure VESTA [29] packages was used. VASPKIT was used as the post-processing tool for VASP code [30].

3. Results and discussion

3.1. Geometrical optimization

The addition of external impurities to the pure TiO₂ leads to structural distortions in the lattice, which eventually influences its electronic, photocatalytic and optical properties [31]. The geometrical distortions are reflected by changes in the local coordination of the atom and the variations in bond length [32]. The relaxed geometry and its distortion of Ti–O, Sc–O, and V–O coordination are shown in Fig. 2, and the lattice parameters and atomic position are given in Table 1. The lattice parameters for pure anatase TiO₂ was obtained after the complete structural relaxation as $a = 3.81 \text{ \AA}$, $b = 3.81 \text{ \AA}$ and $c = 9.70 \text{ \AA}$ are in good agreement with the previously studied experimental [33], and the theoretical values [34]. Ti atom in pure anatase TiO₂ is coordinated to six O atoms surrounding it. Thus formed Ti–O₆ octahedra has four equivalent planar Ti–O (1.95 Å) bonds and two axial Ti–O bonds (2.00 Å). As seen from Fig. 2 (a), the axial O–Ti–O bonds almost form 180° bond angle, whereas the planar O–Ti–O bonds are bent significantly,

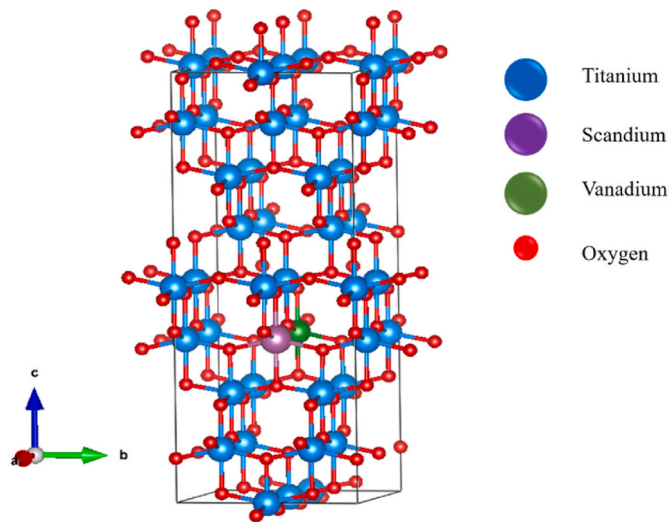


Fig. 1. The relaxed structure of $2 \times 2 \times 2$ supercell of Sc and V co-doped anatase TiO_2 . The Sc, V, Ti, and O atoms are highlighted by purple, green, blue, and red respectively. (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

affecting the electronic structure of TiO_2 .

The incorporation of Sc atom on the Ti site leads increase in bond length due to the elongation of bond, thus Sc–O coordination has the bond length of 2.04 Å and 2.15 Å, in its planer and axial bonds, respectively. Sc doping leads to the distortion of planar oxygen atoms O_{p1} and O_{p2} (Fig. 2(b)) in upward direction, the angle between the planar and axial oxygen atoms in top and bottom are 75.60° and 104.39° respectively, and the angle between the planar oxygen atoms is 93.54°. The addition of V atom on the Ti site leads to the contraction in bond length, the V–O coordination has the bond length of 1.92 Å and 1.93 Å, in its planer and axial bonds respectively. V doping leads to the distortion of planar oxygen atom atoms O_{p1} and O_{p2} (Fig. 2(c)) in downward direction, the angle between the planar and axial oxygen atoms in top and bottom are 78.44° and 101.55° respectively, and the angle between

the planar oxygen atoms is 92.29°.

For the case of Sc and V co-doped structure, the bond length in Sc–O coordination is elongated and the V–O coordination is contracted, and for the Ti–O coordination the bond length is similar to the pure TiO_2 . In the co-doped system, the planar and axial bond lengths are 2.09 Å, 2.11 Å and 1.82 Å, 1.95 Å for Sc–O and V–O bonds respectively. For the co-doped system, in the Sc–O coordination, the angle between the planar and axial oxygen atoms are 76.63°, and 103.36° in top and bottom respectively, the angle between the planar oxygen atoms are 93.07°. For the case of V–O coordination in the co-doped system, the angle between the planar and axial oxygen atoms are 76.89°, and 103.10° in top and bottom respectively, the angle between the planar oxygen atoms are 92.94°. The atomic radius for Sc is 160 p.m. which is higher than Ti (140 p.m.), hence Sc doping increases the cell volume to 1%. The atomic radius for V is 135 p.m., and hence the cell volume is get decreased to 1% by V doping. Similarly, the co-doping of both Sc and V increases its volume by 1%. After the co-doping, the symmetry of the structure was changed from tetragonal $P-4m2$ to orthorhombic $Pmm2$ space group, our study agrees well with the previous study done by Xiong et al. [35]. Table 1 gives the detailed geometrical structure of the system.

3.2. Defect formation energy

The defect formation energy is calculated in order to find the stability of the host compound after introducing the intrinsic and extrinsic defects. The defect formation energy can be calculated by the following formula,

$$\Delta H^{def}(q) = E_{tot}^{def} - E_{tot}^{bulk} \pm \sum_i n_i \mu_i q [E_V + E_F] \quad (1)$$

Where E_{tot}^{def} and E_{tot}^{bulk} denotes the total energy of doped and pure supercell respectively, E_F and E_V are the fermi energy and valence band maximum respectively, n_i indicates the total number of atoms and q is the charge state, μ_i is the chemical potential of the defect species [36]. The chemical potential of the constituent species such as Sc, Ti, and V were calculated from the formation energy of metallic Sc (hcp), Ti (hcp), and V (bcc), respectively.

The calculated defect formation energy for Sc-doped, V-doped and Sc

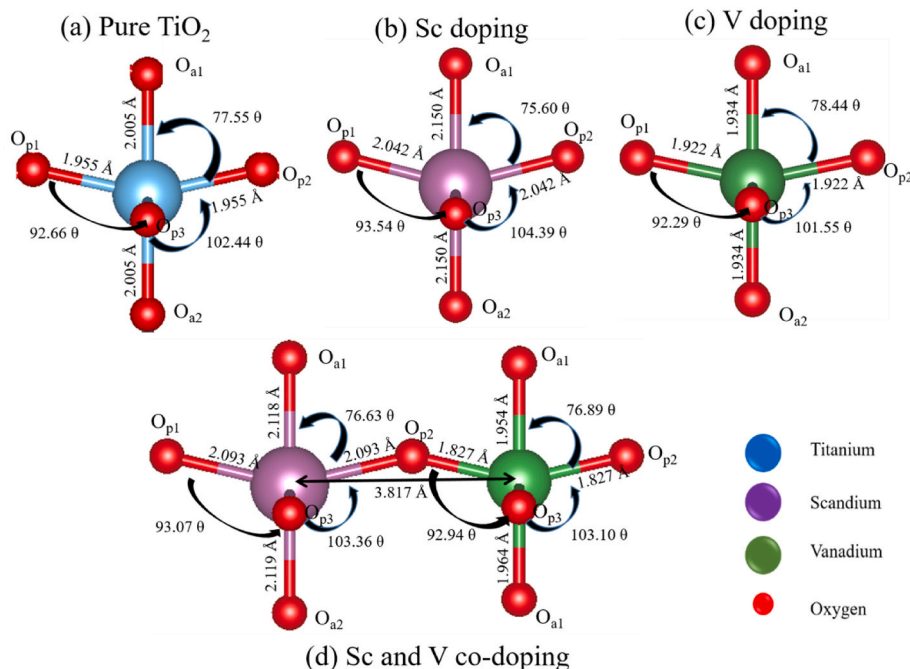


Fig. 2. Geometrical distortion of Ti–O coordination, Sc–O coordination, V–O coordination and Sc and V co-doping.

Table 1

The calculated lattice parameters, cell volume, and bond length for mono doped, and co-doped systems.

Composition	Lattice parameters (Å)				Calculated bond length (Å) and Angle (°)					
					Ti-O		Sc-O		V-O	
	a	b	c	V	Planar	Axial	Planar	Axial	Planar	Axial
Pure	3.819	3.819	9.708	141.58	1.955	2.005	–	–	–	–
Sc -doped	3.827	3.827	9.717	142.31 (1%)	92.66°	102.44°	2.042	2.150	–	–
					1.974	2.045				
					1.906	2.060				
V -doped	3.819	3.819	9.690	141.32 (–1%)	93.34°	104.04°	–	–	1.922	1.934
					1.954	1.993				
					1.977	1.996				
Sc & V codoped	3.817	3.820	9.741	142.03 (1%)	92.40°	101.75°	2.024	2.118	1.827	1.954
					1.953	1.980				
					1.989	2.031				
					1.928	101.27°				
					93.33°	78.59°				

and V co-doped TiO₂ with various charge defects as a function of Fermi energy is shown in Fig. 3. “The configuration energy defined as the difference in the total energy of the supercell with defects and that of the pure supercell, gives an idea about the stability of system after introducing defects. The calculated configuration energy obtained from the optimized total energies was a negative value for Sc in 3+ charge state (–0.0841 eV) and V in 2+ (–0.4775 eV) and 3+ (–0.9584 eV) charge states. This implies that the anatase TiO₂ remains stable when Sc and V are introduced as defects into the TiO₂ lattice.”

The formation energy at zero charge states for Sc doped, V doped, and Sc and V co-doped TiO₂ is 1.079 eV, 2.404 eV, and 1.466 eV, respectively. This implied that Sc doping is energetically favourable. But, the formation energy of Sc–V co-doped TiO₂ is less than that of V doped TiO₂. Hence the Sc doping not only favours charge compensation, but also improves the stability of the co-doped system, which agrees with the experimental observation of similar co-doping done by Murzin et al. [24].

3.3. Electronic structure analysis

The projected density of states (PDOS) for pure anatase TiO₂ is shown in Fig. 4(a). From the hybrid functional calculations, the calculated band gap of pure anatase TiO₂ is about 3.48 eV which is in good agreement with the previous study [37], whereas the experimental bandgap is about 3.2 eV [38]. From the PDOS analysis, it is seen that the valence band is contributed by the O-2p orbitals and the conduction

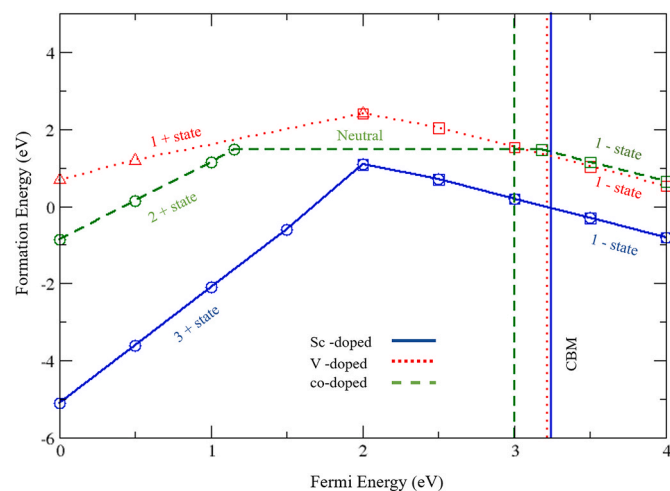


Fig. 3. Calculated defect formation energy with GGA-PBE functionals, for Sc doped, V doped, and Sc, V co-doped TiO₂ at different charge (q) states with respect to the Fermi Energy (E_F).

band is mainly contributed by the Ti-3d orbitals.

Usually Sc has the valence state of 3+, whereas Ti has 4+ valence state. The DOS analysis shows that the Sc doping at the Ti site leads to reduction in the band gap by about 0.24 eV. From the PDOS given in Fig. 4(b), no impurity state contributed by Sc is seen in the forbidden gap. The Sc dopant is supposed to introduce acceptor energy levels close to VBM. This is consistent with the small extra peak of PDOS seen at the E_F (Fig. 4(b)), having the width of about 0.3 eV and mainly made of O-2p orbitals. Sc doping increase in the hole mobile charge carriers which is reflected by the shifting of fermi energy level towards the valence band maximum (VBM).

Generally, V can take up the oxidation states 2+, 3+, 4+ and 5+ in oxides. The PDOS of V-doped TiO₂ is given in Fig. 4(c). It can be seen that the band gap of TiO₂ is reduced by ~0.26 eV due to V doping. Moreover, the E_F moves towards the CBM, and a few additional states appear close to the CBM (0–0.5eV). The Ti-d and V-d states are degenerate in this energy region.

To understand the effect of ambipolar doping of Sc and V into TiO₂, the plots of PDOS of Sc, V co-doped anatase TiO₂ is shown in Fig. 4(d). Unlike the V-doped case, the E_F lies close to VBM in the co-doped system. Interestingly additional states are seen close to the CBM in the energy range of 3–3.5 eV. This isolated DOS region is mainly contributed by V-3d orbitals.

The effect of Sc and V doping on the band features of TiO₂ is schematically shown in Fig. 5. The addition of Sc in co-doping drives the separation of V-3d states that are close to CBM, which reduces the bandgap of TiO₂ from 3.48 eV to 3.0 eV. The band gap obtained for the pure system is 3.48 eV (from hybrid functional calculations) which is over estimated by 0.28 eV than the experimental band gap value (3.2 eV). By including the scissor operation, the band gap value obtained for Sc and V doping are 2.96 eV, and 2.94 eV, respectively, after correcting the over-estimation error. Similarly, a band-gap value of 2.72 eV was obtained for the Sc–V co-doped TiO₂. Thus Sc and V co-doped TiO₂ absorb more visible light, and thus can eventually enhance the photocatalytic activity in the visible region.

3.4. Chemical bonding analysis

To understand the nature of chemical bonding of the mono doped and co-doped system, the charge density (Fig. 6) and the Electron localization function (Fig. 7) analysis were carried out. The contour plots of the electron localization function along the 001 plane for pure, Sc doped, V doped, and Sc, V co-doped anatase TiO₂, is shown in Fig. 7. Even though the electronegativity difference between Ti and O is 1.9, Ti–O bond has significant covalent characteristics in addition to the ionic nature. This is consistent with the previous study by C Sousa et al. [39]. The local structure of TiO₂ is disturbed by Sc doping. It is interesting to note that Sc–O bonds show more covalent character even

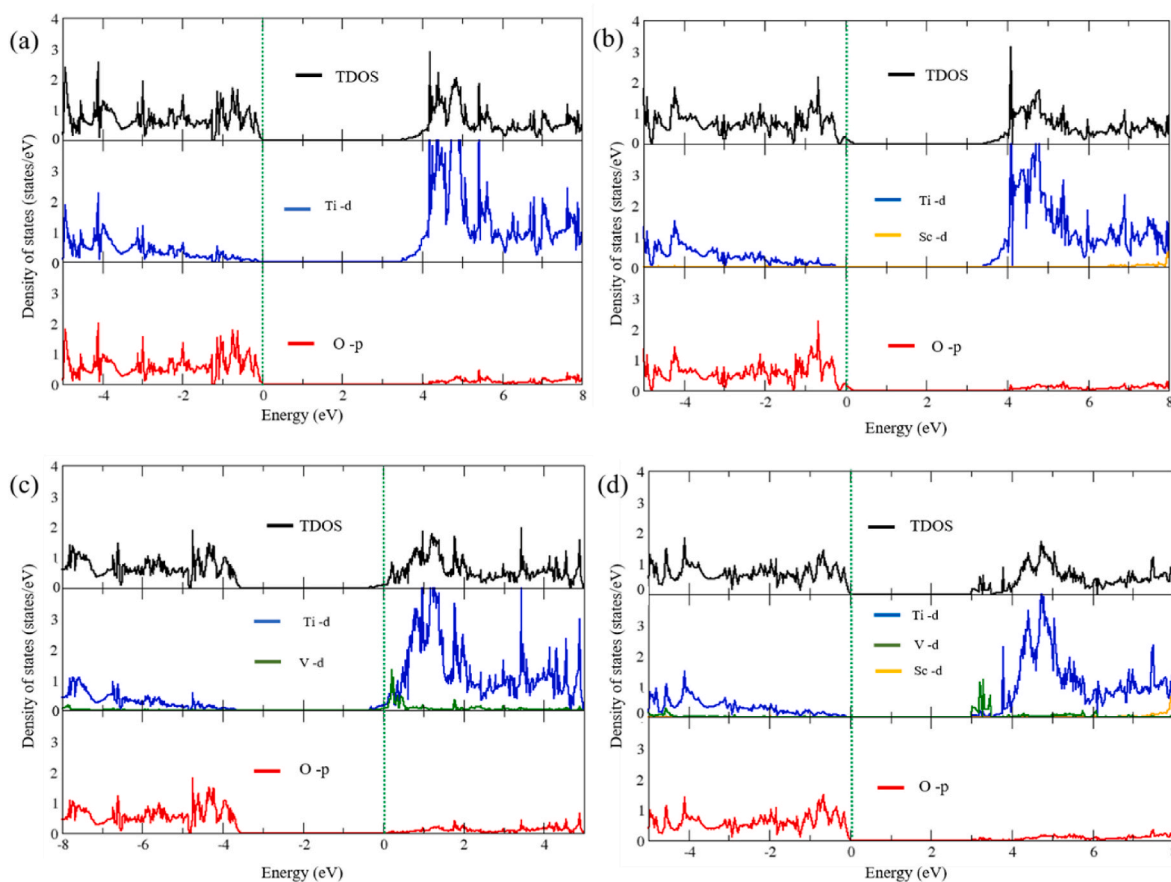


Fig. 4. Projected Density of states plots for (a) pure, (b) Sc doped, (c) V doped, and (d) Sc, V co-doped anatase TiO_2 . The fermi level is shifted to 0 eV, indicated by vertical dotted green lines. (For interpretation of the references to colour in this figure legend, the reader is referred to the Web version of this article.)

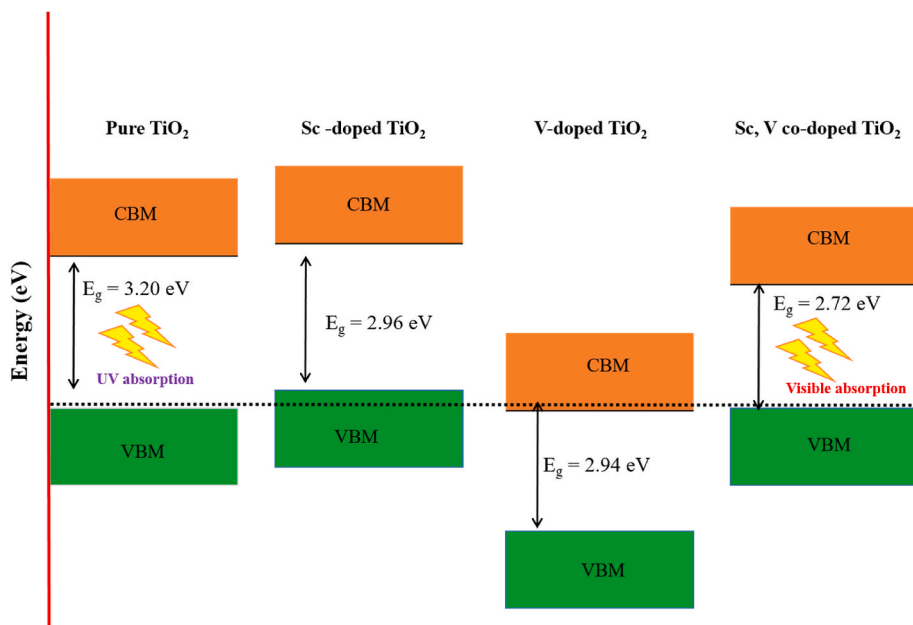


Fig. 5. The schematic representation for band edge positions for pure, monodoped, and co-doped anatase TiO_2 . The fermi level is shifted to 0 eV.

though the electronegativity difference is 2.08 (higher than that of pristine Ti–O). However, the nature of ionic bonding is also non-negligible. Owing to the presence of Sc, the lattice is locally distorted in which the Ti–O bonds undergoing significant changes. The

basal Ti–O bonds exhibit more covalent character and the Sc–O bond lengths are elongated compared to the Ti–O bonds (observed from the charge density and ELF plots in Figs. 6 and 7).

Since the electronegativity bonds between V–O is smaller (1.81) than

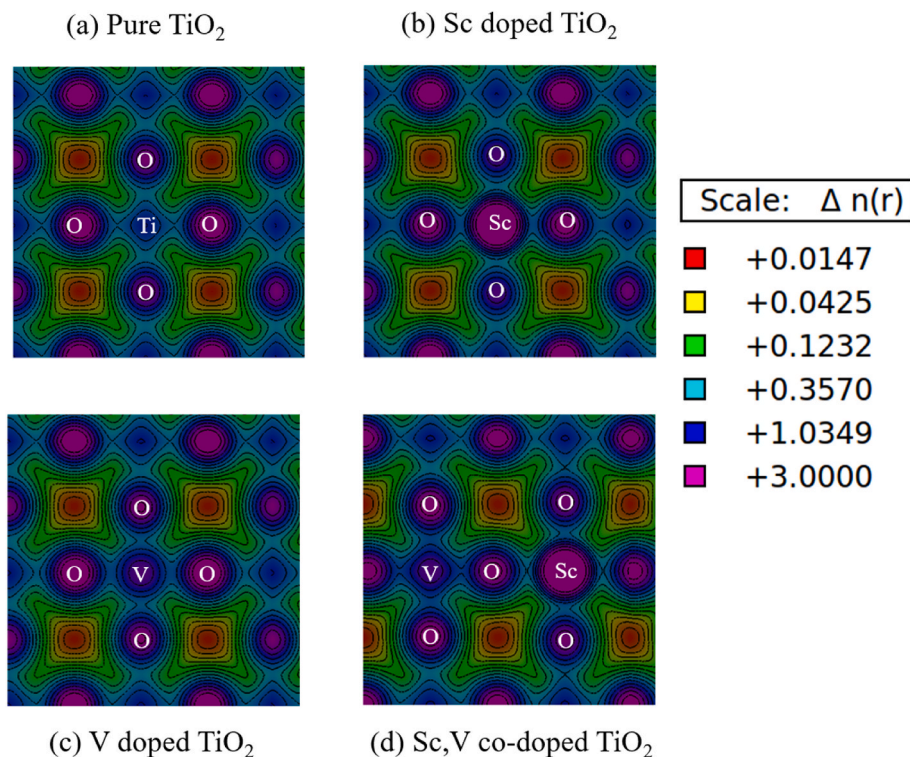


Fig. 6. The Charge density for pure (a), Sc -doped (b), V -doped (c), and Sc, V codoped (d) anatase TiO_2 along (001) plane.

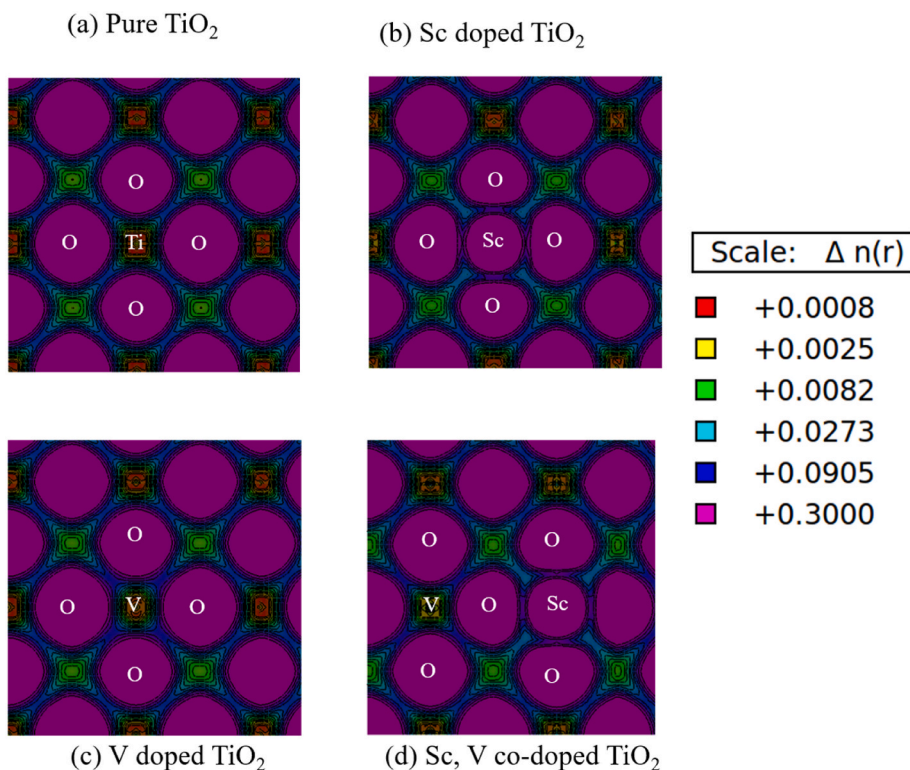


Fig. 7. The Electron localization function for pure (a), Sc -doped (b), V -doped (c), and Sc, V codoped (d) anatase TiO_2 along (001) plane.

that of Ti–O (1.9). The V–O bonds are expected to be more predominantly covalent. More numbers of contours seen between the charge density and ELF plots (Figs. 6 and 7) are consistent with the following assumption. The significant reduction in the V–O bond length (seen along the basal as well as axial planes) compared to the Ti–O bonds,

could be resulting because of the stronger V–O covalent bonds. It is interesting to analyze the effect of ambipolar doping (V as donor and Sc as acceptor) in TiO_2 . The elongation in the Sc–O bond length and the contraction in the V–O bond length seen in the monodoping are also maintained in the co-doping case of TiO_2 .

3.5. Optical properties analysis

The optical properties of the materials are calculated from the frequency dependent complex dielectric function which consists of real [$\epsilon_1(\omega)$] and the imaginary [$\epsilon_2(\omega)$] parts, as stated by the following equation,

$$\epsilon(\omega) = \epsilon_1(\omega) + i\epsilon_2(\omega) \quad (2)$$

The real and imaginary part of the dielectric function shows the capacity of light reflection and absorption of the material respectively [40]. The imaginary [$\epsilon_2(\omega)$] part is calculated from summing over the empty bands, and the real part [$\epsilon_1(\omega)$] is obtained by Kramers-Kronig transformations [$\epsilon_2(\omega)$] [41,42]. The absorption coefficient is calculated from the following equation [43],

$$\alpha(\omega) = \left[2 \left[\sqrt{\epsilon_1^2(\omega) + \epsilon_2^2(\omega)} - \epsilon_1(\omega) \right] \right]^{1/2} \quad (3)$$

Fig. 8 (a) shows the absorption coefficient of the pure TiO₂, Sc doped TiO₂, V doped TiO₂, and Sc, V co-doped TiO₂, which is calculated from equation (3). It can be seen that, the absorption coefficient of pure TiO₂, and the Sc doped TiO₂ is almost similar. The absorption starts from 2.5 eV, corresponding to the UV region. Since V doping introduces electron charge carriers, which brings in metal like nature (Fig. 4(c)), a broad absorption peak is seen at the IR region. Sc and V co-doping leads to improvement in the optical absorption coefficient (Fig. 8(a)). The insert in Fig. 8(a) clearly shows that, the optical absorption starts from 2 eV in the co-doped system. The co-doped system achieves red-shift absorption compared to the pure and mono doped systems. The optical absorption plots of the pure, mono doped, and co-doped system replicates the band structure plots of all the TiO₂ systems, where the band gap is considerably reduced (given in Fig. 1 of the Supplementary Information). Due to Sc and V charge compensated co-doping, the additional energy states formed at the CBM, decreases the electronic transition energy, which eventually improves the light absorption towards the visible light. Fig. 8 (b) shows the imaginary part of dielectric constant as a function of energy. In the energy range from 2.5 eV to 3.5 eV, the pure, mono doped, and co-doped TiO₂ has two absorption peaks at 2.8 eV and 3.4 eV, respectively. From 2 eV the co-doped system shows enhanced optical absorption, compared to mono doped and pure TiO₂. From the optical spectra, we can ensure that, the Sc, V co-doped TiO₂ increases the absorption and also expands the light absorption edge towards the visible region.

4. Conclusion

The structural, electronic, and optical properties of the anatase TiO₂,

which is mono doped with Sc and V, and their co-doping are studied by Density Functional Theory. It may be noted that V-doping in V⁵⁺ charge state becomes a donor, whereas Sc-doping in Sc³⁺ state becomes an acceptor, resulting in charge compensation in the co-doped TiO₂. In addition, this ambipolar impurity combination induces a large local structural distortion in the form of Jahn-Teller distortion, resulting in enhanced stability for the co-doped system. The formation energy for Sc_{Ti} and V_{Ti} in various charge states is calculated as a function of fermi energy. To obtain the bandgap of pure and doped systems accurately, the hybrid functional calculation (HSE06) was used. The band gap obtained for the pure system is 3.48 eV which is over-estimated by 0.28 eV than the experimental band gap value (3.2 eV). From the scissor operator, the band gap value obtained for Sc and V doping are 2.96 eV and 2.94 eV respectively, after including the over-estimation error. Similarly, a band-gap value of 2.72 eV was obtained for the Sc-V co-doped TiO₂. The reduction in the band-gap is resulting due to the additional energy states, formed below the conduction band minimum. This implies that Sc-V co-doping can lead to the absorption in visible energy range, which will further enhance the photo-conversion efficiency of TiO₂.

The excited-state properties like optical absorption and dielectric tensor are calculated for the pure and doped systems. It is found that, the Sc and V co-doping shifts the optical absorption from 2.5 eV to 2.0 eV, which corresponds to the visible region. The bonding characteristics were studied by the electronic charge density and the electron localization function analyses. The V–O bonds are expected to be more predominantly covalent. The Sc–O bonding shows more ionic nature than Ti–O bonding. The elongation in the Sc–O bond and the contraction in the V–O bond seen in the monodoping are maintained in the co-doping case also. Overall we have studied the electronic structure by the effect of Sc/V doping and co-doping on the anatase TiO₂. We found that, Sc/V co-doping can lead to local structural distortion, and enhancement the optical absorption in the visible region.

CRedit authorship contribution statement

S. Muthukrishnan: Data curation, Formal analysis, Writing – original draft, collecting the data, performed the analysis, wrote the paper.
R. Vidya: Data curation, Formal analysis, Writing – original draft, collecting the data, performed the analysis, wrote the paper.
Anja Olafsen Sjøstad: Formal analysis, Conceptualization, designed the analysis.

Declaration of competing interest

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests: Dr R Vidya reports financial support was provided by Anna University.

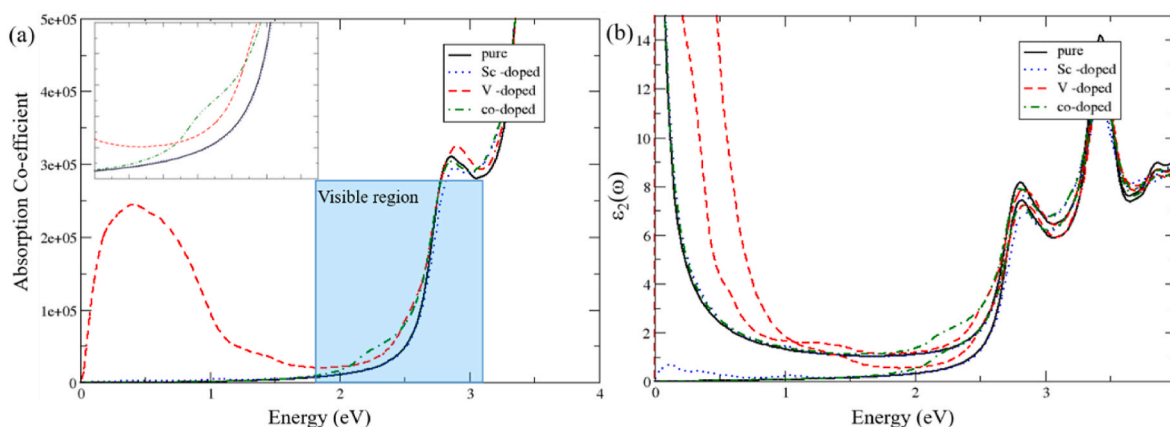


Fig. 8. (a) The absorption coefficient of pure, mono doped, and co-doped TiO₂, the insert indicates shift in optical absorption, (b) the imaginary part of dielectric constant for pure, mono doped, and co-doped TiO₂, as a function of Energy.

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Data availability

Data will be made available on request.

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Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.matchemphys.2023.127467>.

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